

D0-D0bar Mixing in an exclusive approach

Fu-Sheng Yu

Lanzhou University

Collaboration with: H-n. Li, C-D. Lu, Q. Qin

in progress

CHARM @ WSU

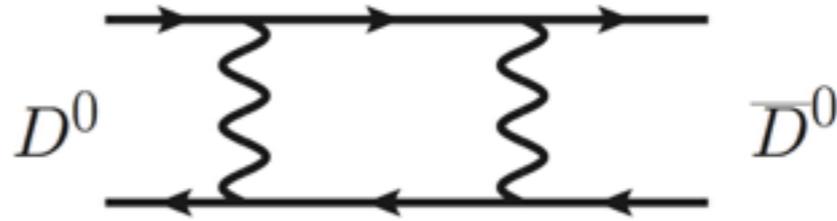
20.05.2015

Outline

A.Kagan's talk tomorrow

- Motivation
 - DDbar mixing misunderstood in theory
- D-Dbar Mixing in an exclusive approach
 - Large flavor SU(3) breaking effect in factorization-assisted topological-amplitude approach
- Conclusions

$D^0 - \bar{D}^0$ Mixing



- The time evolution

$$i \frac{\partial}{\partial t} \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix} = \left(\mathbf{M} - \frac{i}{2} \mathbf{\Gamma} \right) \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix}$$

- Mass eigenstates in terms of weak eigenstates

$$|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle$$

- Mass difference and Width difference

$$x \equiv \frac{\Delta m}{\Gamma} = \frac{m_1 - m_2}{\Gamma}$$

$$y \equiv \frac{\Delta \Gamma}{2\Gamma} = \frac{\Gamma_1 - \Gamma_2}{2\Gamma}$$

Experiment vs. Theory

- Current world average results

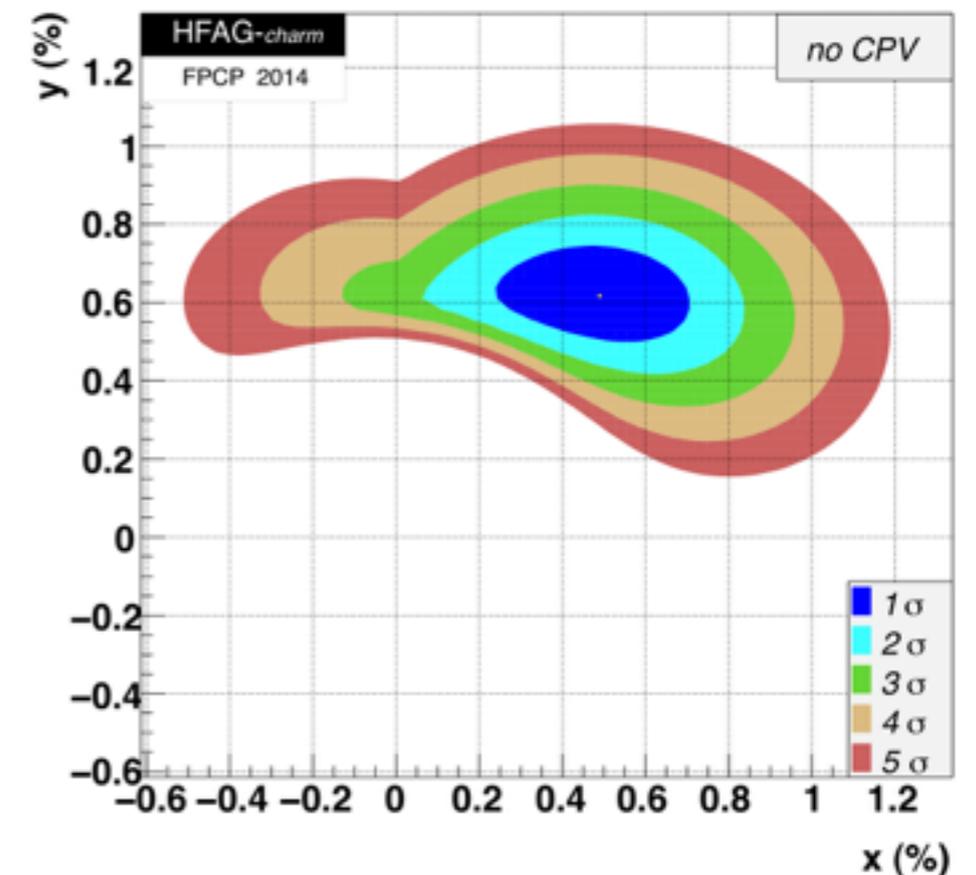
[HFAG, 2014]

- * If CP is conserved

$$x = (0.49^{+0.14}_{-0.15})\%, \quad y = (0.62 \pm 0.08)\%$$

- * If CP violation is allowed

$$x = (0.41^{+0.14}_{-0.15})\%, \quad y = (0.63^{+0.07}_{-0.08})\%$$



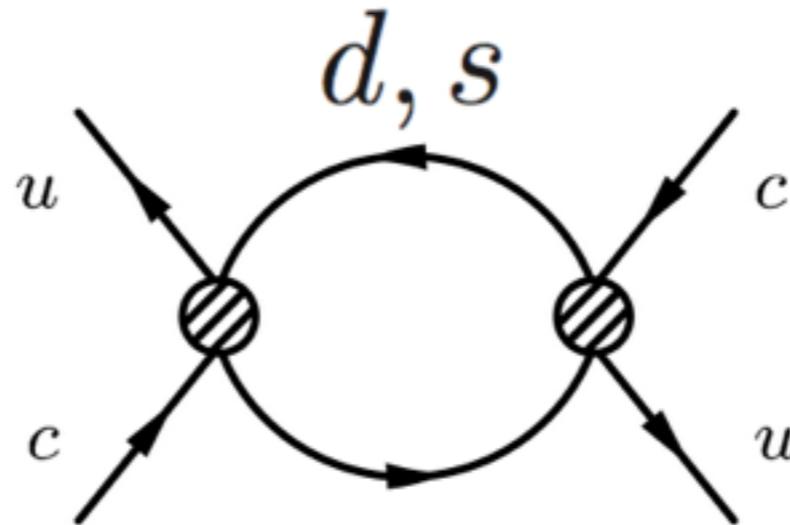
- So far, **quantitatively** theoretical calculation can just reach to the order of $x \sim y \sim 0.1\%$

[Bigi et al, 00'; Lenz et al, 10'; Cheng et al, 10']

Inclusive approach:

quark level

Heavy Quark Expansion



$$x \sim (m_s/m_c)^4$$

$$y \sim (m_s/m_c)^6$$

$$x \sim 10^{-6}$$

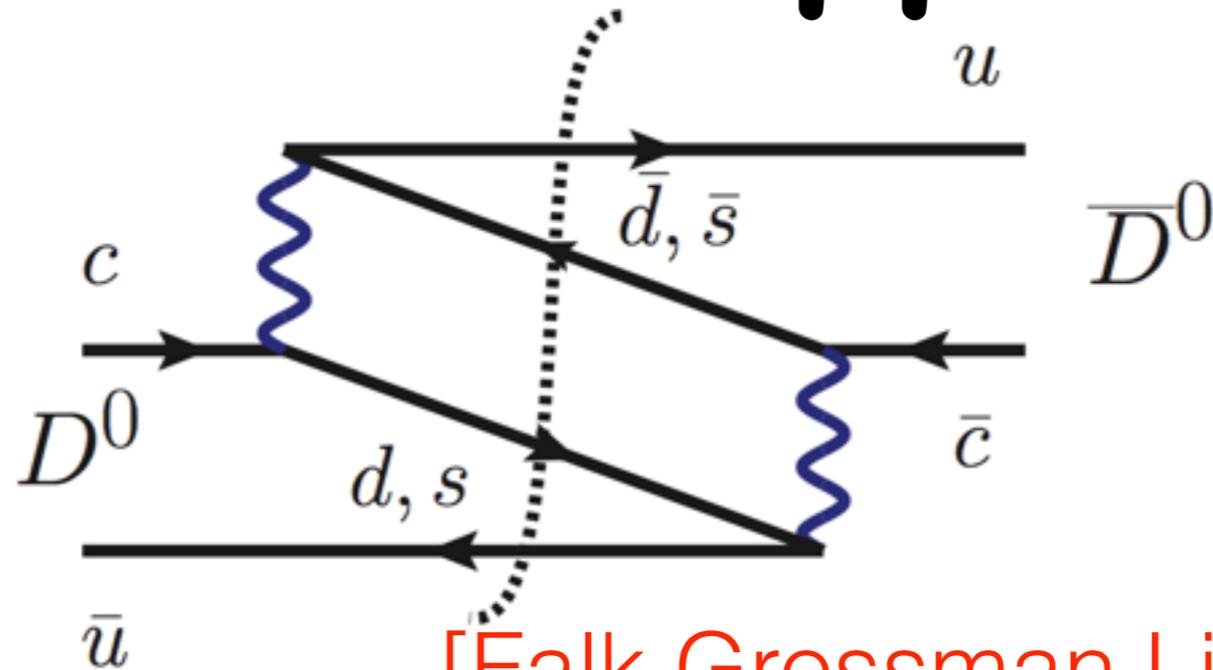
$$y \lesssim 0.9 \times 10^{-5}$$

[Bobrowski, Lenz, 09']

[Lenz, et al, 10']

Short-distance contributions are small

Exclusive Approach



[Falk, Grossman, Ligeti, Petrov, 02']

$$y = \frac{1}{2\Gamma} \sum_n \rho_n \eta_{\text{CP}}(n) (\langle D^0 | H_w | n \rangle \langle \bar{n} | H_w | D^0 \rangle + \langle D^0 | H_w | \bar{n} \rangle \langle n | H_w | D^0 \rangle)$$

$$= \sum_n \eta_{\text{CKM}}(n) \eta_{\text{CP}}(n) \cos \delta_n \sqrt{\text{Br}(D^0 \rightarrow n) \text{Br}(D^0 \rightarrow \bar{n})},$$

Sum up all the intermediate states

⇒ To predict Branching fractions well

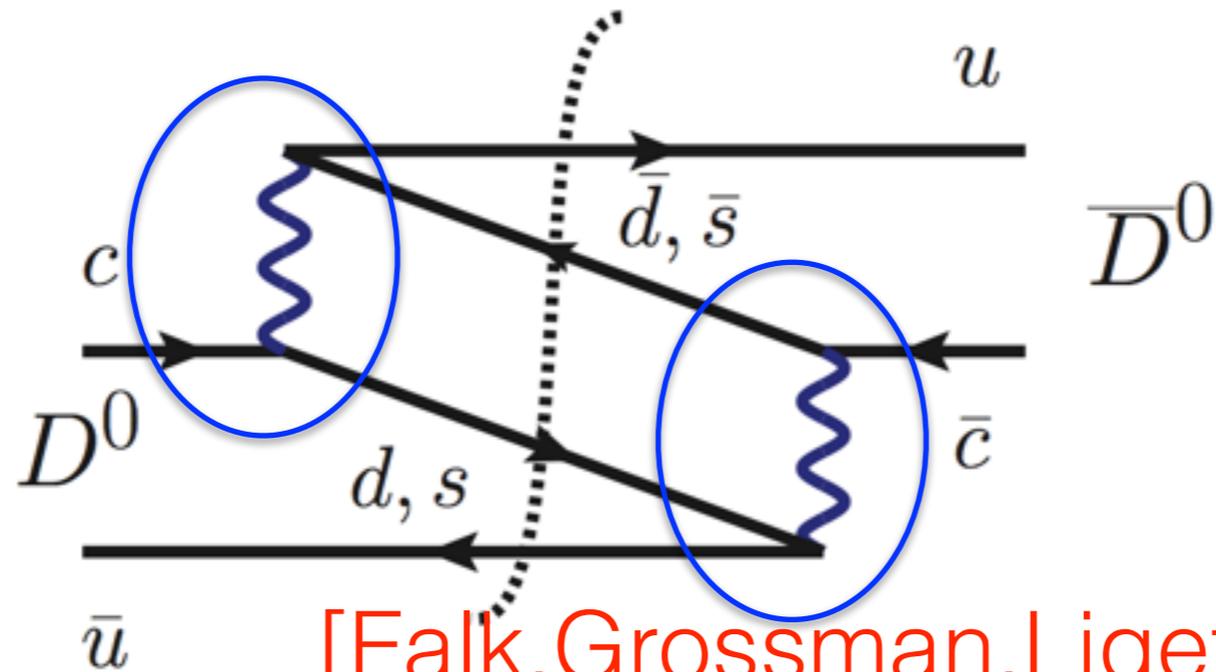
Exclusive Approach

Sum up all the intermediate states

$$y = \frac{1}{2\Gamma} \sum_n \rho_n \eta_{\text{CP}}(n) (\langle D^0 | H_w | n \rangle \langle \bar{n} | H_w | D^0 \rangle + \langle D^0 | H_w | \bar{n} \rangle \langle n | H_w | D^0 \rangle)$$
$$= \sum_n \eta_{\text{CKM}}(n) \eta_{\text{CP}}(n) \cos \delta_n \sqrt{\text{Br}(D^0 \rightarrow n) \text{Br}(D^0 \rightarrow \bar{n})},$$

- **SU(3) breaking effect from phase space**
[Falk, Grossman, Ligeti, Petrov, 02'] $\pi\pi\pi\pi$ v.s. $KKKK$
- **Topological diagrammatic approach**
[Cheng, Chiang, 10']
- **U-spin breaking effect** [Gronau, Rosner, 12']

Naive Expectation



$$x, y \sim \sin^2 \theta_C \times [SU(3) \text{ breaking}]^2$$

$$0.2^2 \times 0.3^2 \sim \boxed{0.4\%}$$

Exp: $x = (0.49_{-0.15}^{+0.14})\%$, $y = (0.62 \pm 0.08)\%$

- Well description of **SU(3) breaking** is a **central issue**

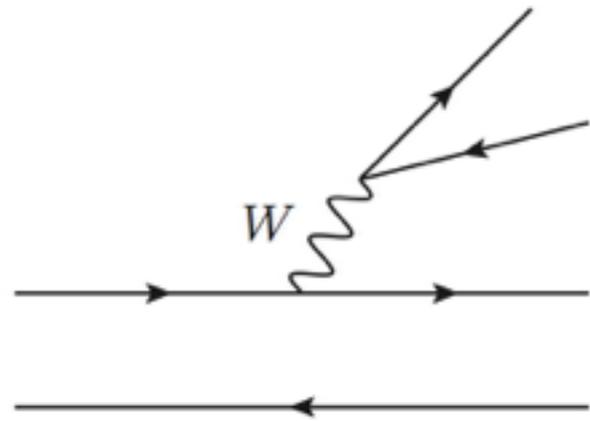
[Paul's talk yesterday]

Problem on dynamics

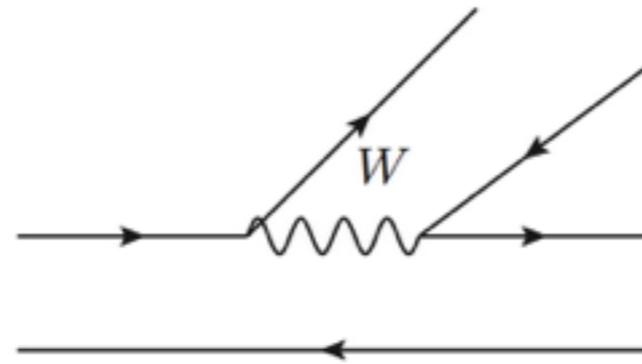
- $m_c \sim 1.3\text{GeV}$
 - Neither heavy enough for heavy quark expansion, $1/mc$
 - Nor light enough for chiral perturbation theory
- QCD-inspired methods do not work: HQET, QCDF, PQCD, SCET.

Large nonperturbative contributions

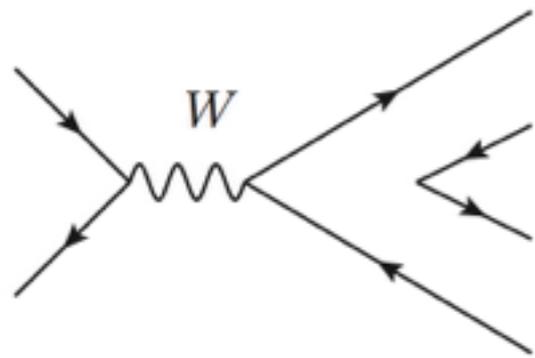
Topological-Amplitude Approach



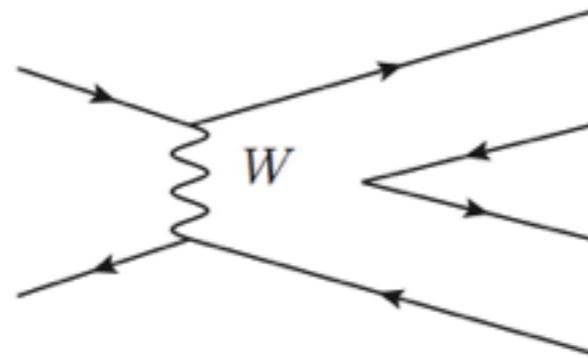
(a) T



(b) C



(c) A



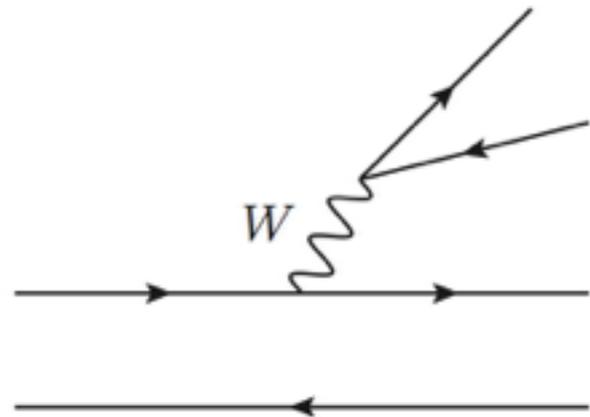
(d) E

[Bhattacharya, Rosner, 08', 10']

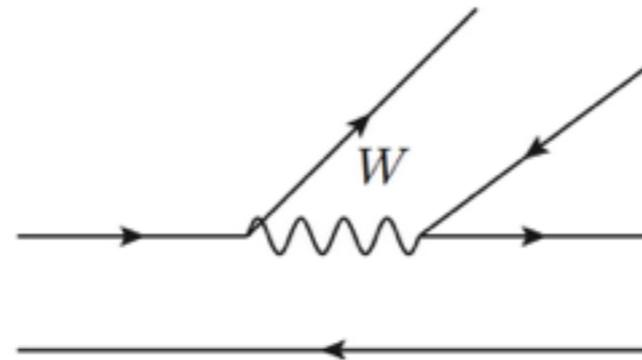
[Cheng, Chiang, 10']

[Schacht's talk tomorrow]

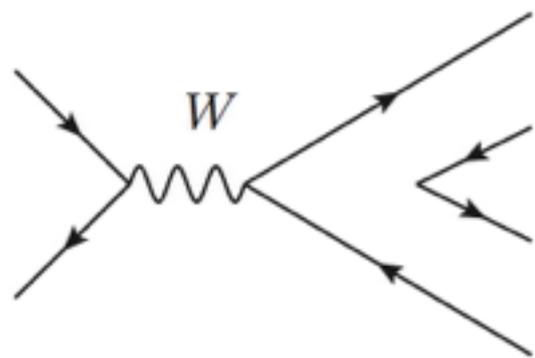
Factorization-Assisted Topological-Amplitude Approach



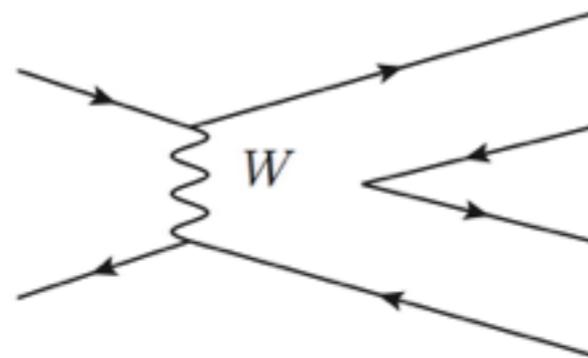
(a) T



(b) C



(c) A



(d) E

[Bhattacharya, Rosner, 08', 10']

[Cheng, Chiang, 10']

[Schacht's talk tomorrow]

- Calculate each topological amplitude in factorization:

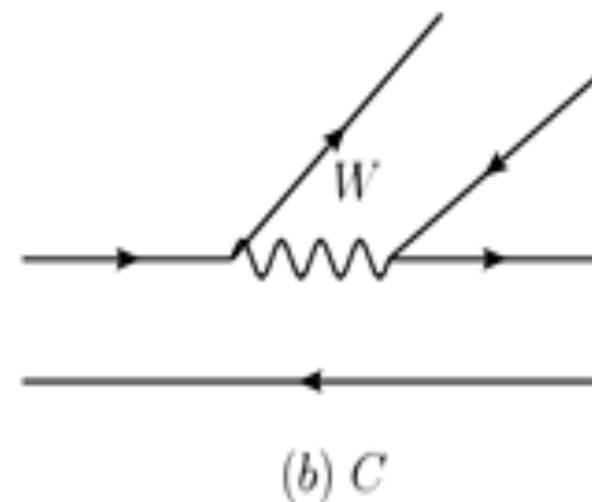
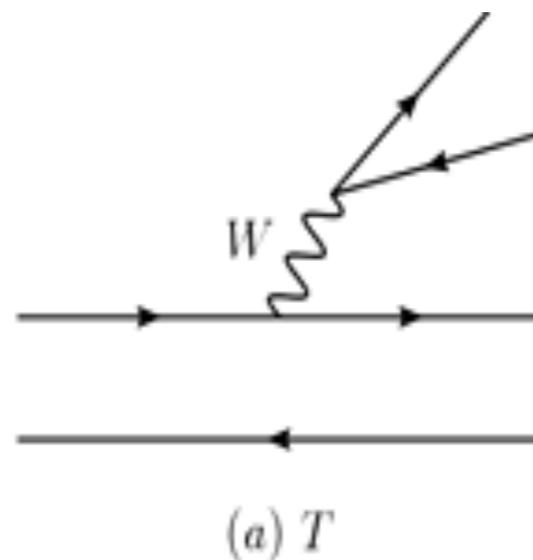
- Short-distance dynamics: Wilson coefficients

- Long-distance dynamics: hadronic matrix elements

[Li, Lu, FSY, 1203.3120]

Emission Amplitudes

- Color-favored Tree (T)
- Color-suppressed (C)



$$\langle P_1 P_2 | \mathcal{H}_{eff} | D \rangle_{T,C} = \frac{G_F}{\sqrt{2}} V_{CKM} a_{1,2}(\mu) f_{P_2} (m_D^2 - m_{P_1}^2) F_0^{DP_1}(m_{P_2}^2)$$

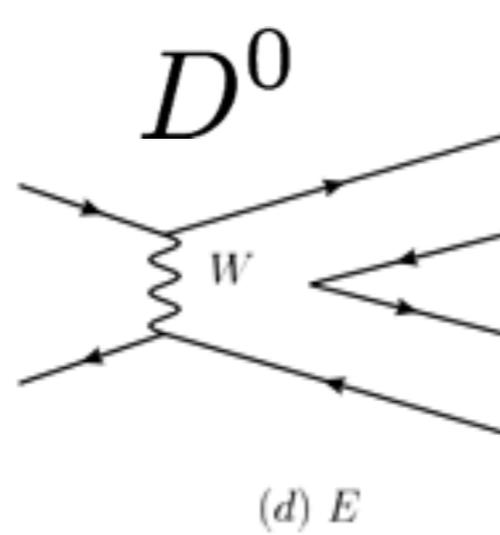
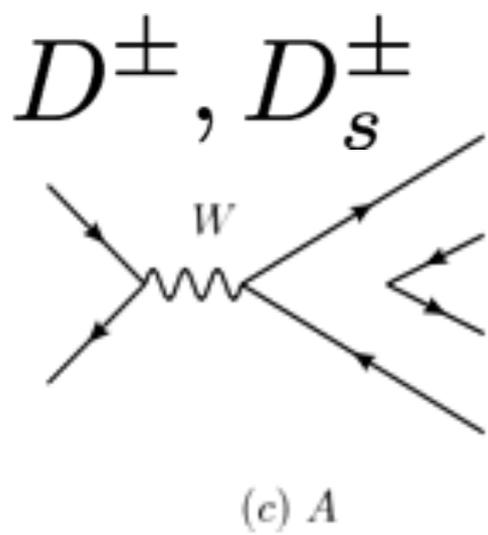
$$a_1(\mu) = C_2(\mu) + \frac{C_1(\mu)}{N_c}, \quad \text{Non-factorizable}$$

$$a_2(\mu) = C_1(\mu) + C_2(\mu) \left[\frac{1}{N_c} + \chi_{nf} e^{i\phi} \right],$$

$$\mu = \sqrt{\Lambda m_D (1 - r_2^2)}, \quad r_2 = m_{P_2}^2 / m_D^2$$

**SU(3)
breaking
effects**

[Li, Lu, FSY, 1203.3120]



[Li, Lu, FSY, 12']

W-annihilation (A) W-exchange (E)

$$\langle P_1 P_2 | \mathcal{H}_{\text{eff}} | D \rangle_{E,A} = \frac{G_F}{\sqrt{2}} V_{\text{CKM}} b_{q,s}^{E,A}(\mu) f_D m_D^2 \left(\frac{f_{P_1} f_{P_2}}{f_\pi^2} \right)$$

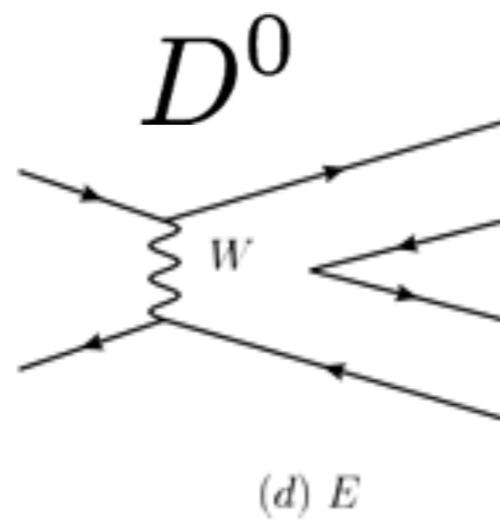
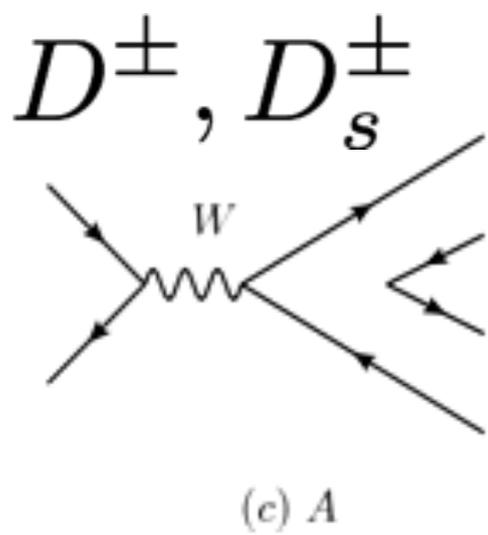
Dominated by **non-factorizable** contribution

A: $b_{q,s}^A(\mu) = C_1(\mu) \chi_{q,s}^A e^{i\phi_{q,s}^A}$

E: $b_{q,s}^E(\mu) = C_2(\mu) \chi_{q,s}^E e^{i\phi_{q,s}^E}$

**Factorization-Assisted
Topological Approach**

nonperturbative
contributions



[Li, Lu, FSY, 12']

W-annihilation (A) W-exchange (E)

$$\langle P_1 P_2 | \mathcal{H}_{\text{eff}} | D \rangle_{E,A} = \frac{G_F}{\sqrt{2}} V_{\text{CKM}} b_{q,s}^{E,A}(\mu) f_D m_D^2 \left(\frac{f_{P_1} f_{P_2}}{f_\pi^2} \right)$$

Dominated by **non-factorizable** contribution

A: $b_{q,s}^A(\mu) = C_1(\mu) \chi_{q,s}^A e^{i\phi_{q,s}^A}$

E: $b_{q,s}^E(\mu) = C_2(\mu) \chi_{q,s}^E e^{i\phi_{q,s}^E}$

**SU(3)
breaking
effects**

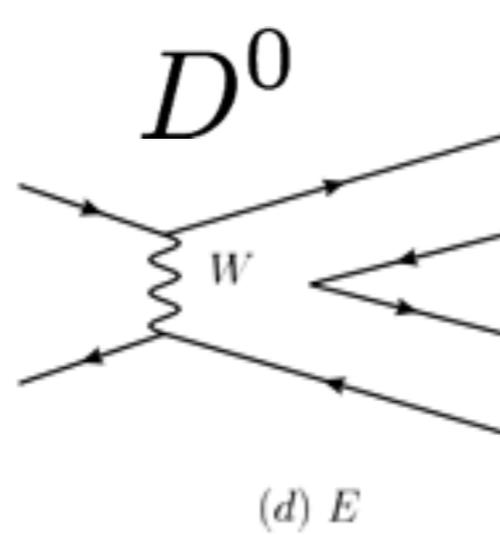
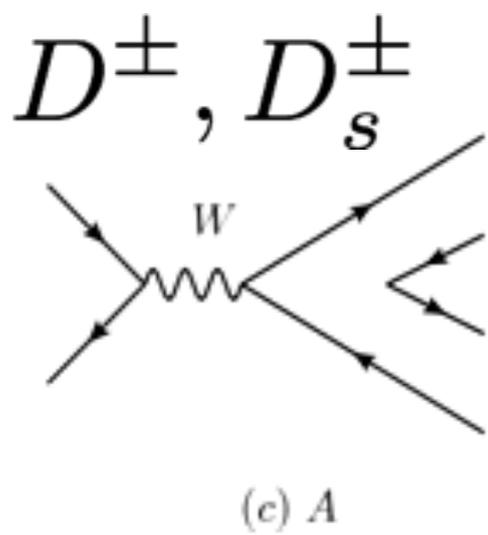
$$\chi^A \sim \chi^E$$

$$\chi_q^E = 0.11, \quad \chi_q^A = 0.12,$$

$$\chi_s^E = 0.18, \quad \chi_s^A = 0.17.$$

$$\chi_q \neq \chi_s$$

subscripts: quark pairs produced from vacuum



[Li, Lu, FSY, 12']

W-annihilation (A) W-exchange (E)

$$\langle P_1 P_2 | \mathcal{H}_{\text{eff}} | D \rangle_{E,A} = \frac{G_F}{\sqrt{2}} V_{\text{CKM}} b_{q,s}^{E,A}(\mu) f_D m_D^2 \left(\frac{f_{P_1} f_{P_2}}{f_\pi^2} \right)$$

Dominated by **non-factorizable** contribution

A: $b_{q,s}^A(\mu) = C_1(\mu) \chi_{q,s}^A e^{i\phi_{q,s}^A + S_\pi}$

E: $b_{q,s}^E(\mu) = C_2(\mu) \chi_{q,s}^E e^{i\phi_{q,s}^E + S_\pi}$

**SU(3)
breaking
effects**

Glauber strong phase for pions

pion => Goldstone boson?
qqbar bound state?

$$\chi_q \neq \chi_s$$

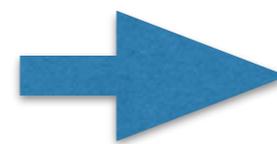
Much and precise experimental data



To extract nonperturbative parameters from data

PP: 12 parameters for 28 data

PV: 14 parameters for 33 data



more dynamics

more predictive

Singly Cabibbo-suppressed modes agree well with experiments

More SU(3) breaking effects included

($\times 10^{-3}$)

Modes	Br(FSI)	Br(diagram)	Br(pole)	Br(exp)	Br(this work)
$D^0 \rightarrow \pi^+ \pi^-$	1.59	2.24 ± 0.10	2.2 ± 0.5	1.45 ± 0.05	1.43 
$D^0 \rightarrow K^+ K^-$	4.56	1.92 ± 0.08	3.0 ± 0.8	4.07 ± 0.10	4.19 
$D^0 \rightarrow K^0 \bar{K}^0$	0.93	0	0.3 ± 0.1	0.320 ± 0.038	0.36
$D^0 \rightarrow \pi^0 \pi^0$	1.16	1.35 ± 0.05	0.8 ± 0.2	0.81 ± 0.05	0.57
$D^0 \rightarrow \pi^0 \eta$	0.58	0.75 ± 0.02	1.1 ± 0.3	0.68 ± 0.07	0.94
$D^0 \rightarrow \pi^0 \eta'$	1.7	0.74 ± 0.02	0.6 ± 0.2	0.91 ± 0.13	0.65
$D^0 \rightarrow \eta \eta$	1.0	1.44 ± 0.08	1.3 ± 0.4	1.67 ± 0.18	1.48
$D^0 \rightarrow \eta \eta'$	2.2	1.19 ± 0.07	1.1 ± 0.1	1.05 ± 0.26	1.54
$D^+ \rightarrow \pi^+ \pi^0$	1.7	0.88 ± 0.10	1.0 ± 0.5	1.18 ± 0.07	0.89
$D^+ \rightarrow K^+ \bar{K}^0$	8.6	5.46 ± 0.53	8.4 ± 1.6	6.12 ± 0.22	5.95
$D^+ \rightarrow \pi^+ \eta$	3.6	1.48 ± 0.26	1.6 ± 1.0	3.54 ± 0.21	3.39
$D^+ \rightarrow \pi^+ \eta'$	7.9	3.70 ± 0.37	5.5 ± 0.8	4.68 ± 0.29	4.58
$D_S^+ \rightarrow \pi^0 K^+$	1.6	0.86 ± 0.09	0.5 ± 0.2	0.62 ± 0.23	0.67
$D_S^+ \rightarrow \pi^+ K^0$	4.3	2.73 ± 0.26	2.8 ± 0.6	2.52 ± 0.27	2.21
$D_S^+ \rightarrow K^+ \eta$	2.7	0.78 ± 0.09	0.8 ± 0.5	1.76 ± 0.36	1.00
$D_S^+ \rightarrow K^+ \eta'$	5.2	1.07 ± 0.17	1.4 ± 0.4	1.8 ± 0.5	1.92

Back to our topic in the end

D-Dbar Mixing in an exclusive approach

$$\begin{aligned} y &= \frac{1}{2\Gamma} \sum_n \rho_n \eta_{\text{CP}}(n) (\langle D^0 | H_w | n \rangle \langle \bar{n} | H_w | D^0 \rangle + \langle D^0 | H_w | \bar{n} \rangle \langle n | H_w | D^0 \rangle) \\ &= \sum_n \eta_{\text{CKM}}(n) \eta_{\text{CP}}(n) \cos \delta_n \sqrt{\text{Br}(D^0 \rightarrow n) \text{Br}(D^0 \rightarrow \bar{n})} , \end{aligned}$$

$$CP|M_1 M_2\rangle = \eta_{\text{CP}}(M_1) \eta_{\text{CP}}(M_2) (-1)^L |M_1 M_2\rangle = \eta_{\text{CP}}(M_1 M_2) |M_1 M_2\rangle$$

\mathcal{YPP}

vanish in the $SU(3)$ symmetry limit

$$\mathcal{B}(\pi^+\pi^-) + \mathcal{B}(K^+K^-) - 2 \cos \delta_{K^+\pi^-} \sqrt{\mathcal{B}(K^-\pi^+) \mathcal{B}(K^+\pi^-)}$$

$$+ \mathcal{B}(\pi^0\pi^0) + \mathcal{B}(K^0\bar{K}^0) - 2 \cos \delta_{K^0\pi^0} \sqrt{\mathcal{B}(\bar{K}^0\pi^0) \mathcal{B}(K^0\pi^0)}$$

$$+ \mathcal{B}(\pi^0\eta) + \mathcal{B}(\pi^0\eta') + \mathcal{B}(\eta\eta) + \mathcal{B}(\eta\eta')$$

$$- 2 \cos \delta_{K^0\eta} \sqrt{\mathcal{B}(\bar{K}^0\eta) \mathcal{B}(K^0\eta)} - 2 \cos \delta_{K^0\eta'} \sqrt{\mathcal{B}(\bar{K}^0\eta') \mathcal{B}(K^0\eta')}$$

$\mathcal{Y}PV$

$$\begin{aligned}
& Br(\pi^0\rho^0) + Br(\pi^0\omega) + Br(\pi^0\phi) + Br(\eta\omega) + Br(\eta'\omega) + Br(\eta\phi) + Br(\eta\rho^0) + Br(\eta'\rho^0) \\
& - 2 \cos \delta_{K^{*-}\pi^+} \sqrt{Br(K^{*-}\pi^+)Br(K^{*+}\pi^-)} - 2 \cos \delta_{K^{*0}\pi^0} \sqrt{Br(K^{*0}\pi^0)Br(\bar{K}^{*0}\pi^0)} \\
& - 2 \cos \delta_{K^-\rho^+} \sqrt{Br(K^-\rho^+)Br(K^+\rho^-)} - 2 \cos \delta_{K^0\rho^0} \sqrt{Br(K^0\rho^0)Br(\bar{K}^0\rho^0)} \\
& - 2 \cos \delta_{K^{*0}\eta} \sqrt{Br(K^{*0}\eta)Br(\bar{K}^{*0}\eta)} - 2 \cos \delta_{K^{*0}\eta'} \sqrt{Br(K^{*0}\eta')Br(\bar{K}^{*0}\eta')} \\
& - 2 \cos \delta_{K^0\omega} \sqrt{Br(K^0\omega)Br(\bar{K}^0\omega)} - 2 \cos \delta_{K^0\phi} \sqrt{Br(K^0\phi)Br(\bar{K}^0\phi)} \\
& + 2 \cos \delta_{K^+K^{*-}} \sqrt{Br(K^+K^{*-})Br(K^-K^{*+})} + 2 \cos \delta_{K^0\bar{K}^{*0}} \sqrt{Br(K^0\bar{K}^{*0})Br(\bar{K}^0K^{*0})} \\
& + 2 \cos \delta_{\pi^+\rho^-} \sqrt{Br(\pi^+\rho^-)Br(\pi^-\rho^+)}
\end{aligned}$$

More decay modes

[Qin, Li, Lu, FSY, 14']

Our results

$$y_{PP} = 0.08\%$$

$$y_{PV} = 0.29\%$$

preliminary

$$y_{PP+PV} = 0.37\%$$

- Experimental data

$$\text{Exp: } y_D = (0.62 \pm 0.08)\% \quad [\text{HFAG}]$$

Outlook 1: uncertainties to be studied

Outlook 2: contributions from other modes

x from dispersion relation

$$\text{Re}f(s) = \frac{1}{\pi} P \int_0^\infty \frac{ds'}{s' - s} \text{Im}f(s')$$

In the heavy quark limit

$$\Delta m = -\frac{1}{2\pi} P \int_{2m_\pi}^\infty dE \left[\frac{\Delta\Gamma(E)}{E - m_D} + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{E}\right) \right]$$

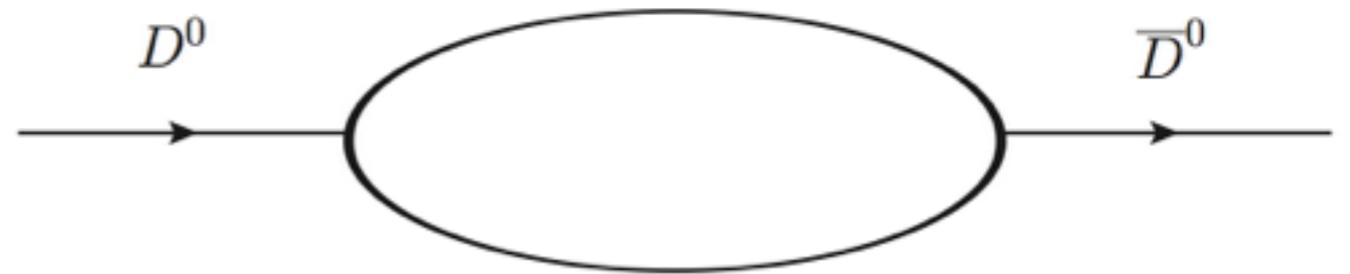
[Falk, Grossman, Ligeti, Nir, Petrov, 2004]

x from dispersion relation

$$\operatorname{Re} f(s) = \frac{1}{\pi} P \int_0^\infty \frac{ds'}{s' - s} \operatorname{Im} f(s')$$

$$\Delta m = 2M_{12} = 2\operatorname{Re}[\Pi(m_D^2)]$$

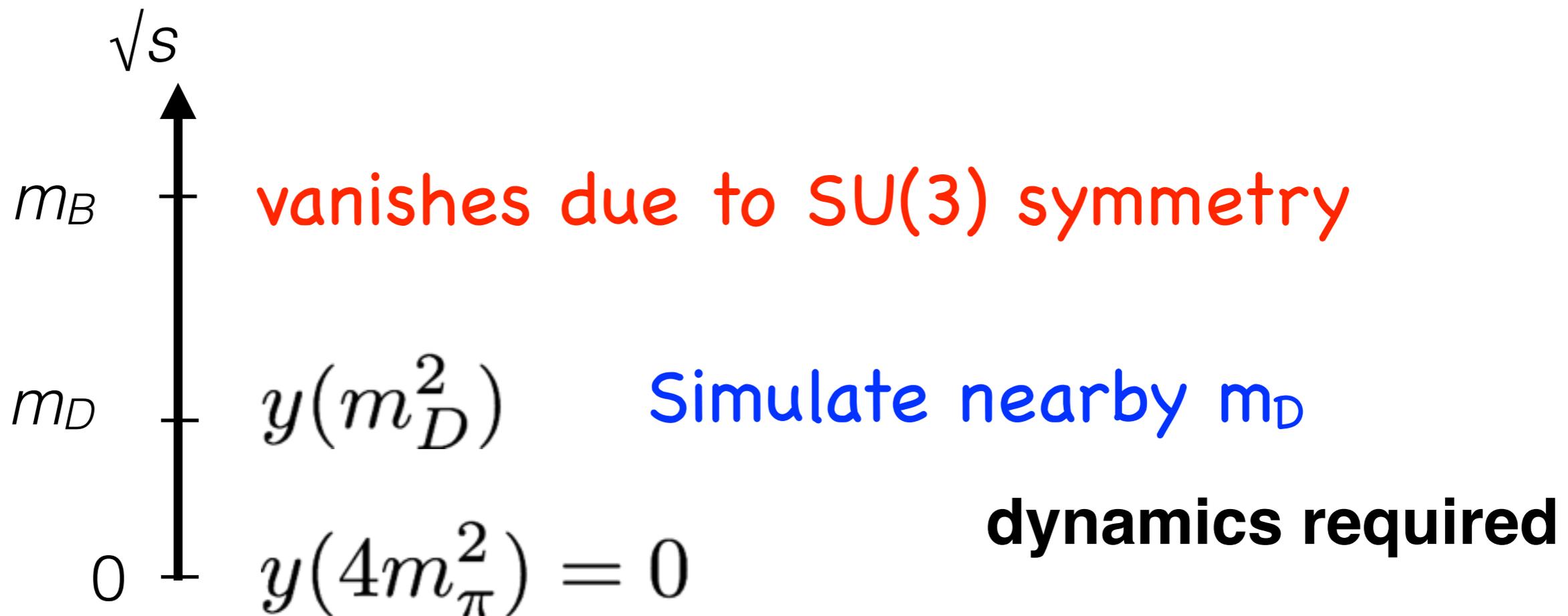
$$\Delta\Gamma = 2\Gamma_{12} = 2\operatorname{Im}[\Pi(m_D^2)]$$



$$\Delta m(m_D^2) = \frac{2}{\pi} P \int_{(2m_\pi^2)^2}^\infty \frac{\Delta\Gamma(s)}{s - m_D^2} ds$$

x from dispersion relation

$$\Delta m(m_D^2) = \frac{2}{\pi} P \int_{(2m_\pi^2)^2}^{\infty} \frac{\Delta\Gamma(s)}{s - m_D^2} ds$$



Summary

- Two-body D meson decays are well studied in Factorization-Assisted Topological-amplitude approach.
SU(3) breaking effects are well described.
- D-Dbar mixing parameter γ can be understood in an exclusive approach.
- More efforts for experimentalist are required to measure VV, PA and PS modes
- More efforts on dynamics are required for mixing parameter χ

Thank you for your attention!

Backup

$$D^0 \rightarrow \pi^+ \pi^- \quad v.s. \quad D^0 \rightarrow K^+ K^-$$

$$\begin{aligned} \mathcal{A}(D^0 \rightarrow \pi^+ \pi^-) &= \frac{G_F}{\sqrt{2}} \lambda_d (T^{\pi\pi} + E^{\pi\pi}) \\ &= \frac{G_F}{\sqrt{2}} V_{cd}^* V_{ud} \left[a_1(\mu) (m_D^2 - m_\pi^2) f_\pi F_0^{D\pi}(m_\pi^2) + C_2(\mu) e^{i(\phi_q^E + 2S_\pi)} \chi_q^E f_D m_D^2 \right] \end{aligned}$$

Glauber phase

Main difference

$$\begin{aligned} \mathcal{A}(D^0 \rightarrow K^+ K^-) &= \frac{G_F}{\sqrt{2}} \lambda_s (T^{KK} + E^{KK}) \\ &= \frac{G_F}{\sqrt{2}} V_{cs}^* V_{us} \left[a_1(\mu) (m_D^2 - m_K^2) f_K F_0^{DK}(m_K^2) + C_2(\mu) e^{i\phi_q^E} \chi_q^E f_D m_D^2 \frac{f_K^2}{f_\pi^2} \right] \end{aligned}$$

Modes	Br(FSI)	Br(diagram)	Br(pole)	Br(exp)	Br(this work)
$D^0 \rightarrow \pi^+ \pi^-$	1.59	2.24 ± 0.10	2.2 ± 0.5	1.45 ± 0.05	1.43 ←
$D^0 \rightarrow K^+ K^-$	4.56	1.92 ± 0.08	3.0 ± 0.8	4.07 ± 0.10	4.19 ←

$$T^{\pi\pi} = 2.73, \quad E^{\pi\pi} = 0.82 e^{-i142^\circ},$$

$$T^{KK} = 3.65, \quad E^{KK} = 1.2 e^{-i85^\circ},$$

Result of CP asymmetries

- Difference of CPV in $D \rightarrow KK$ and $D \rightarrow \pi\pi$

$$\Delta A_{CP}^{\text{dir}} = A_{CP}(D^0 \rightarrow K^- K^+) - A_{CP}(D^0 \rightarrow \pi^- \pi^+)$$

- Our prediction:

$$\Delta A_{CP} = (-0.57 \sim -1.87) \times 10^{-3}$$

[Li, Lu, FSU, 12']

- After our prediction, the world average value is lowered down by the LHCb results

Exp: $\Delta A_{CP} = (-2.53 \pm 1.04) \times 10^{-3}$ [HFAG2014]